

---

# Intensity and pulse shape effect on the spectral and angular distribution of nonlinear Thomson scattering\*

**Yuelin Li and Stephen V. Milton**

*Advanced Photon Source, Argonne National Laboratory*

*Argonne, IL 60565*

\*Supported by the U. S. Department of Energy, Office of Basic Energy Sciences  
Contract No. W-31-109-ENG-38

# Background and related



**Electron dynamics and radiation in high intensity laser fields**

**Ponderomotive scattering**

**Violent acceleration**

**Harmonic and broadband radiation**

....

**Laser acceleration of electrons**

**MeV and GeV electron acceleration**

.....

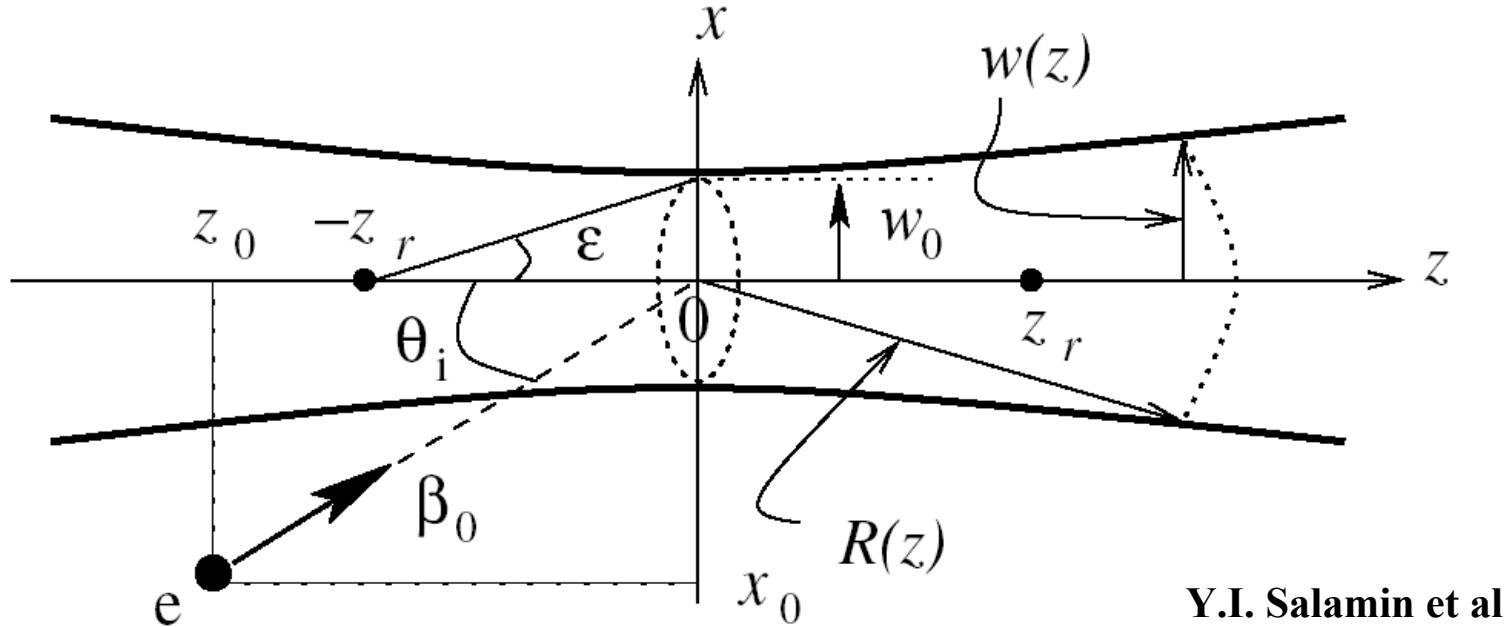
**Ultrafast X-ray radiation generation**

**Nonlinear Thomson scattering into X-ray**

**Bremsstrahlung X-ray**

....

# Geometry



Y.I. Salamin et al

**Only electrons initially at rest and on axis are considered,  
 $x_0=0$ ,  $\beta_0=0$ ,  $\lambda=0.8 \mu\text{m}$ ,  $w_0=10 \mu\text{m}$ ,  $Z_r=400 \mu\text{m}$ .**

## Field: angular spectrum representation

$$E_x = \frac{E_0}{4\epsilon^2} \left( I_1 + \frac{x^2 + y^2}{k_0 r^3} I_2 + \frac{y^2}{r^2} I_3 \right),$$

$$E_y = \frac{E_0}{4\epsilon^2} \frac{xy}{k_0 r^3} (k_0 r I_3 - 2 I_2),$$

$$E_z = \frac{E_0}{4\epsilon^2} \frac{x}{r} I_4,$$

$$B_x = \frac{E_y}{c},$$

$$B_y = \frac{E_0}{4c\epsilon^2} \left( I_1 + \frac{y^2 - x^2}{k_0 r^3} I_2 + \frac{x^2}{r^2} I_3 \right),$$

$$B_z = \frac{E_0}{4c\epsilon^2} \frac{y}{r} I_4$$

$$I_1 = \int_0^1 e^{-b^2/4\epsilon^2} (1 + \sqrt{1 - b^2}) \sin(\phi_b) J_0(k_0 r b) b \, db,$$

$$I_2 = \int_0^1 e^{-b^2/4\epsilon^2} \frac{\sin(\phi_b)}{\sqrt{1 - b^2}} J_1(k_0 r b) b^2 \, db,$$

$$I_3 = \int_0^1 e^{-b^2/4\epsilon^2} \frac{\sin(\phi_b)}{\sqrt{1 - b^2}} J_0(k_0 r b) b^3 \, db,$$

$$I_4 = \int_0^1 e^{-b^2/4\epsilon^2} (1 + \frac{1}{\sqrt{1 - b^2}}) \cos(\phi_b) J_1(k_0 r b) b^2 \, db,$$

$$\phi_b = \omega_0 t - k_0 z \sqrt{1 - b^2} + \phi,$$

$$\epsilon = \frac{1}{k_0 w_0},$$

$$r = \sqrt{x^2 + y^2}.$$

The dynamics is solved numerically

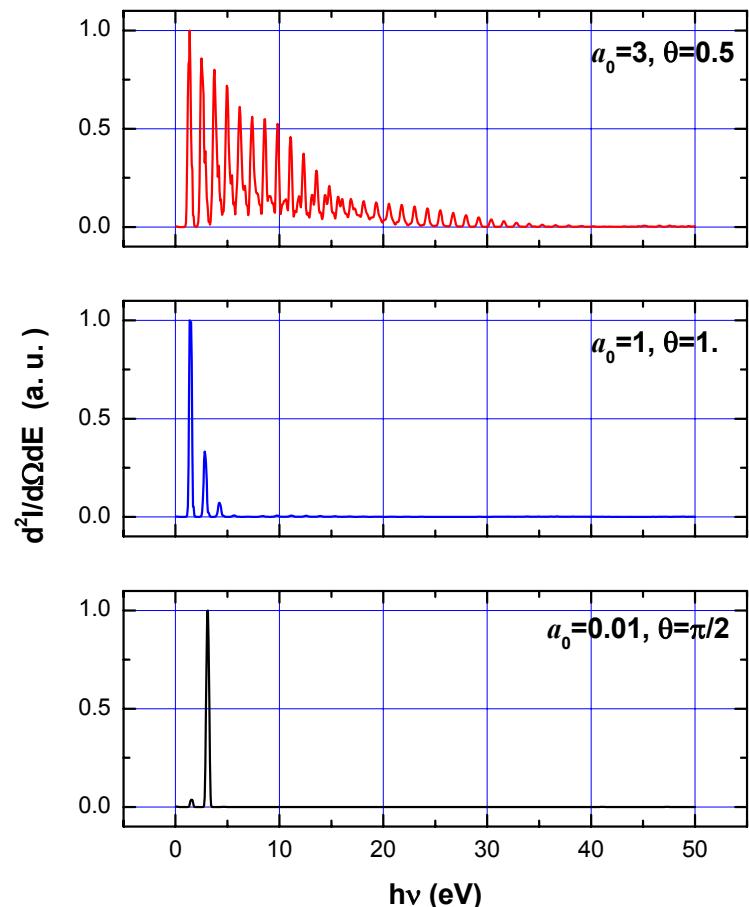
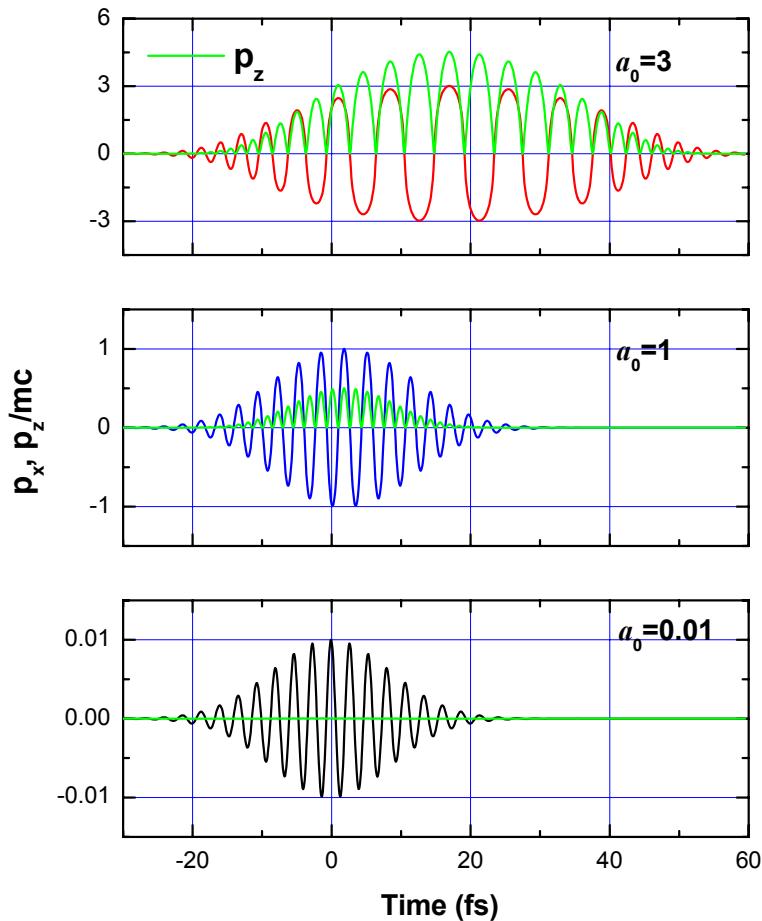
$$\vec{r} = \int_0^t \frac{\vec{p}}{\gamma} dt, \quad \vec{p} = \vec{p}_0 + \int_0^t \dot{\vec{p}} dt, \quad \dot{\vec{p}} = e(\vec{E} + \frac{\vec{p}}{\gamma m} \times \vec{B}).$$

Energy spectrum calculation by Lienard-Wiechert potentials

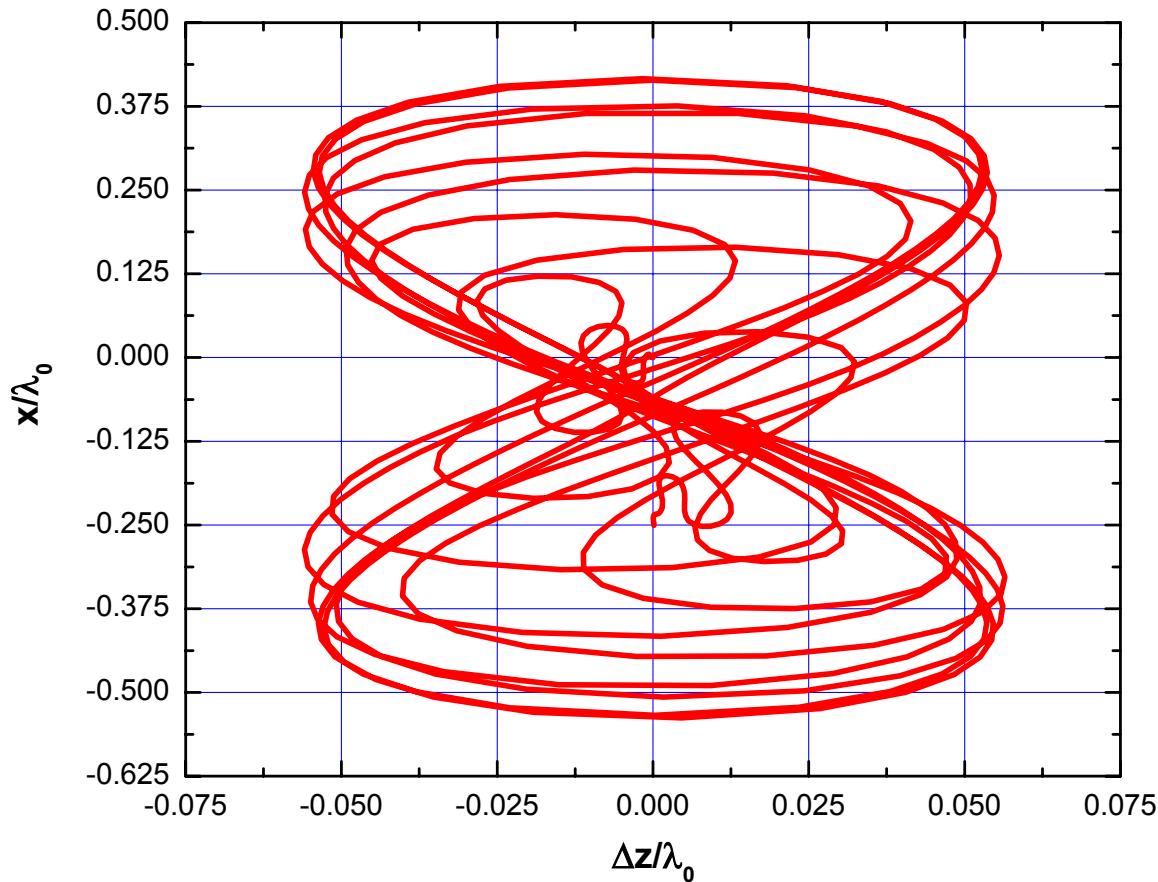
$$\frac{d^2 I}{d\omega d\Omega} = \frac{e^2 \omega^2}{4\pi^2 c} \left| \int \vec{n} \times (\vec{n} \times \vec{\beta}) e^{i\omega(t - \vec{n} \cdot \vec{r}/c)} dt \right|^2.$$

# Single electron case

## Momentum and spectra



# Single electron trajectory, R-frame



# Single electron: Spectrum details

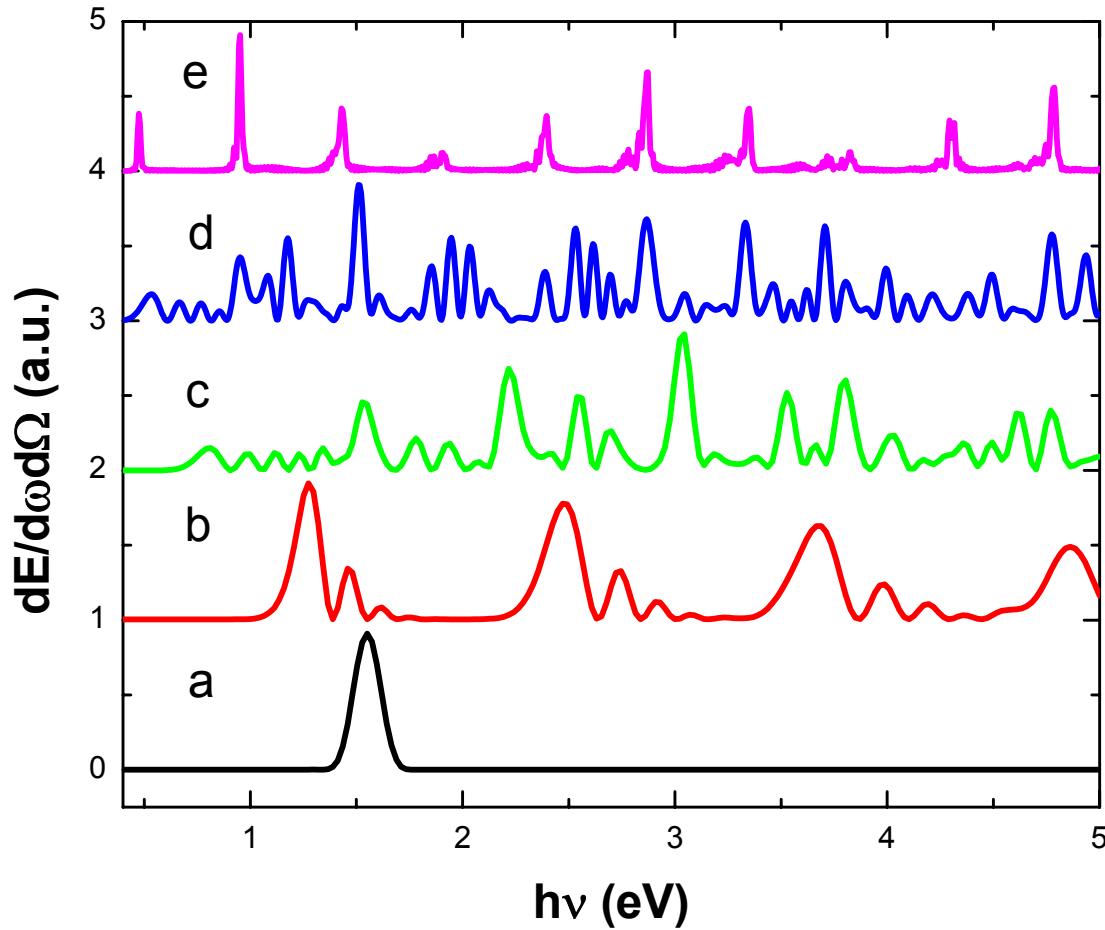


Fig. 1. Spectra of single electron irradiated by (a)-(d) 20-fs laser pulse with  $a_0=3$  for  $\phi=40^\circ$  at  $\theta=0, 30, 60$ , and  $90^\circ$ , and (e) with quasi-flat-top 80-fs pulse with  $a_0=3$  for  $\phi=40^\circ$  at  $\theta=90^\circ$ .

# Single electron: Angular distribution

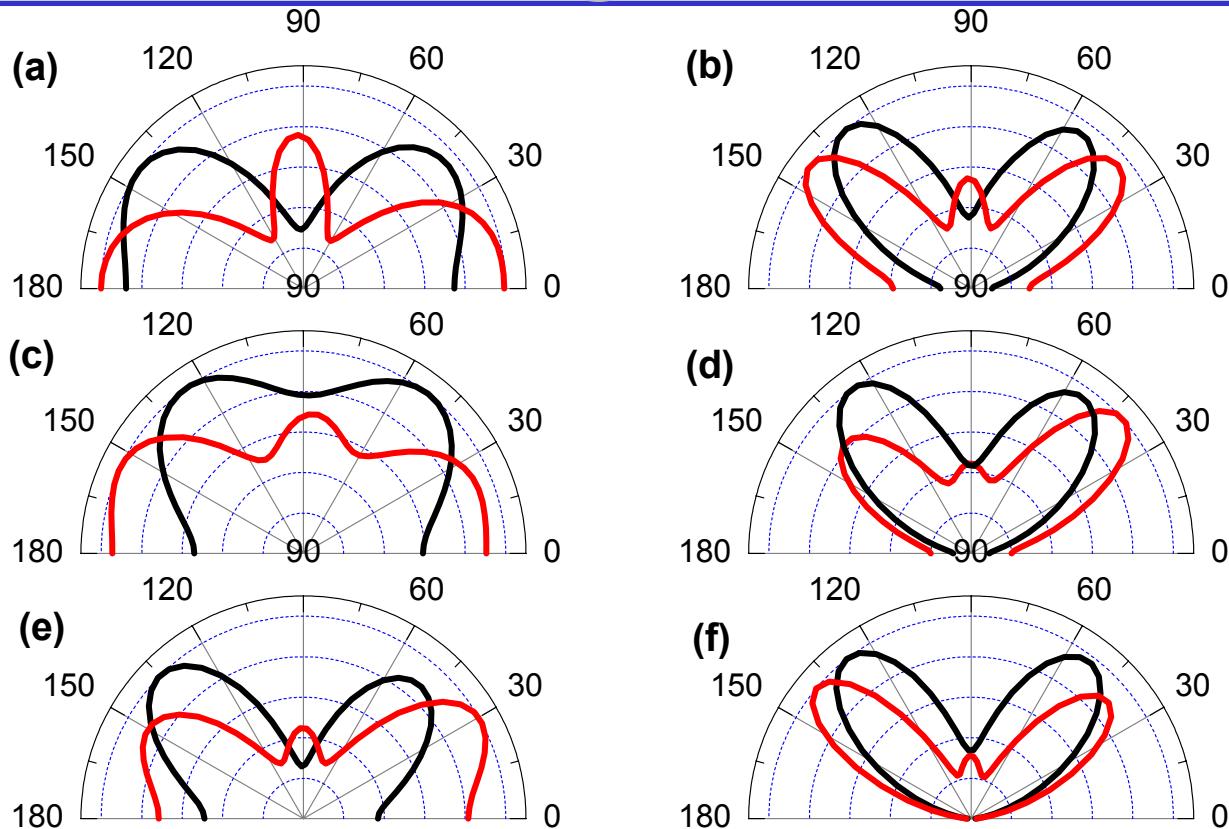
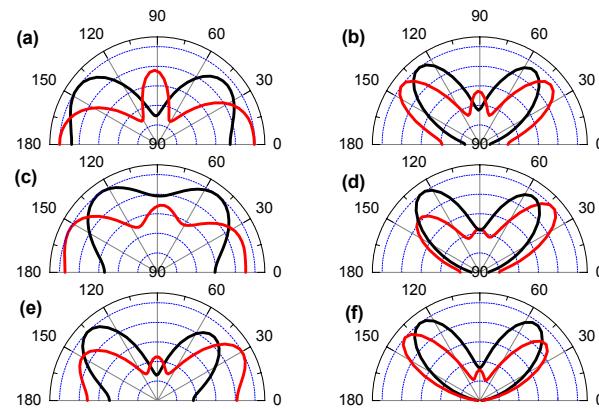
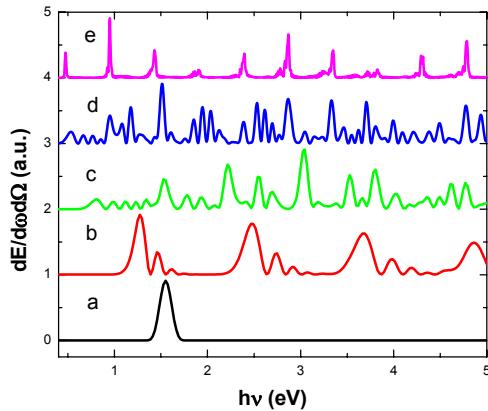


Fig. 2. Radiation pattern of the second (solid) and third (dashed) harmonics at  $\theta=90$  (left panels) and  $60^\circ$  (right panels) for (a)-(b) 80-fs quasi-flat-top pulse  $a_0=3$ , (c)-(d) 20-fs pulse with  $a_0=3$  and (e-f) 20-fs pulse with  $a_0=2$ .

# Single electron: conclusion



## Observations

1. Red shift with intensity and polar angle, explained by theory

$$\omega_n = \frac{n\omega_0}{1 + (a/2)^2(1 - \cos\theta)}$$

2. Spectral broadening, modulation, and over lapping are pulse- shape and intensity- dependent and not discussed previously.
3. Angular distribution is also intensity-and pulse shape-dependent

# Collective effects



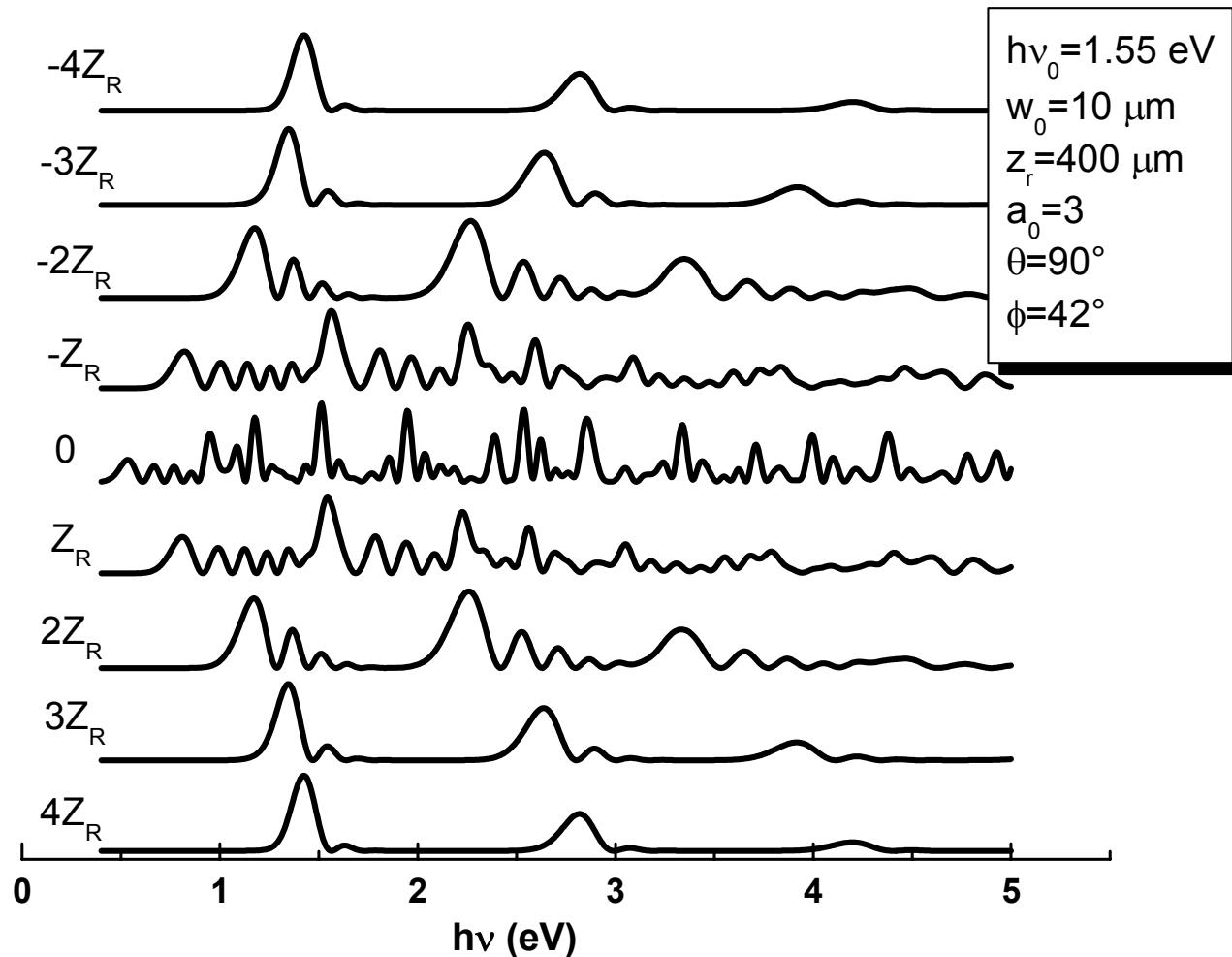
For a collection of electrons, the radiation is calculated as

$$\frac{d^2 I}{d\omega d\Omega} = \frac{e^2 \omega^2}{4\pi^2 c} \left| \int \sum_j \vec{n} \times (\vec{n} \times \vec{\beta}_j) e^{i\omega(t_j - \vec{n} \cdot \vec{r}_j / c)} dt \right|^2.$$

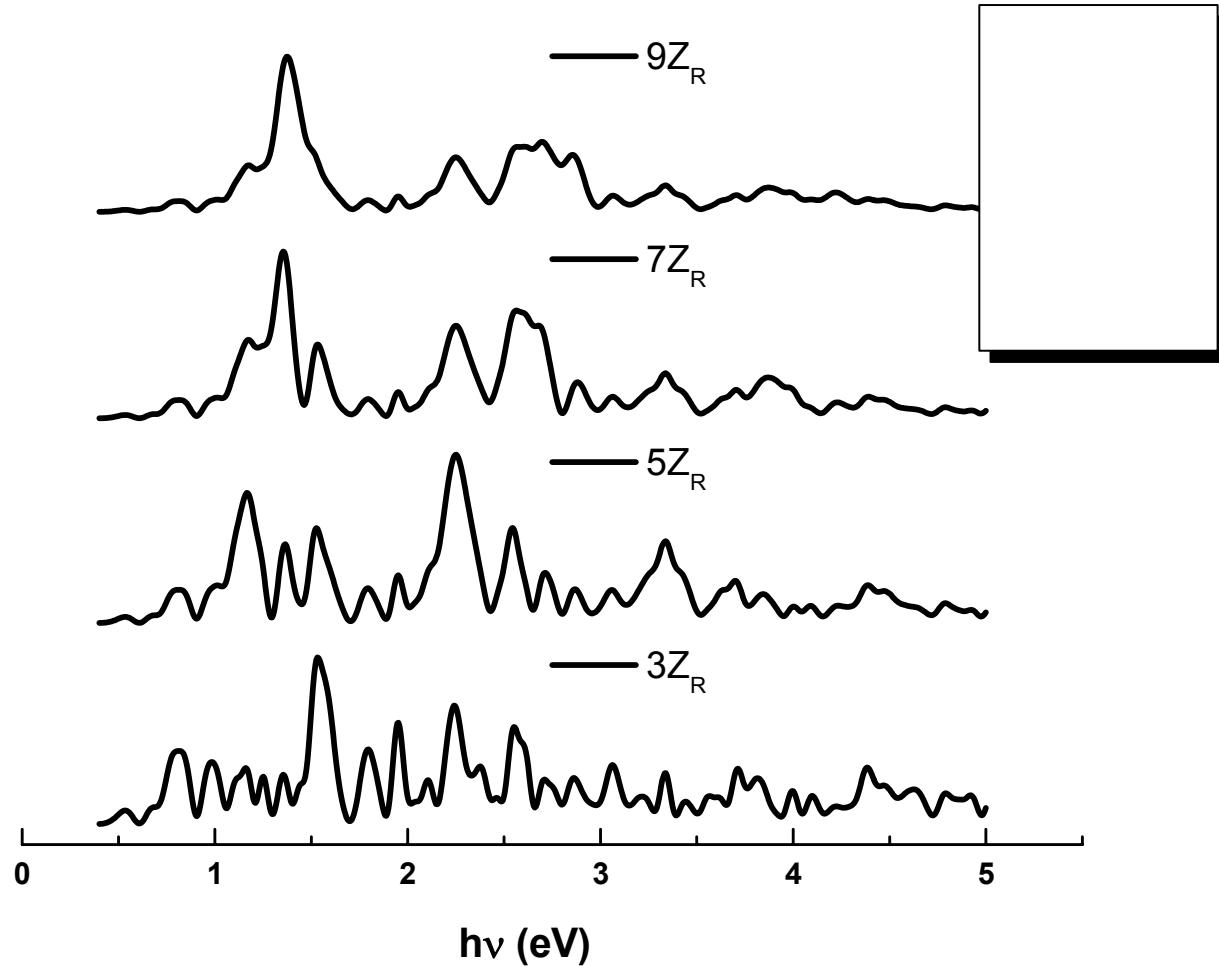
When the positions of the electrons are totally random, the only terms survive the integration are the individual terms of each electrons, hence the radiation for a collection of electrons can be approximated as

$$\frac{d^2 I}{d\omega d\Omega} = \frac{e^2 \omega^2}{4\pi^2 c} \sum_j \left| \int \vec{n} \times (\vec{n} \times \vec{\beta}_j) e^{i\omega(t_j - \vec{n} \cdot \vec{r}_j / c)} dt \right|^2.$$

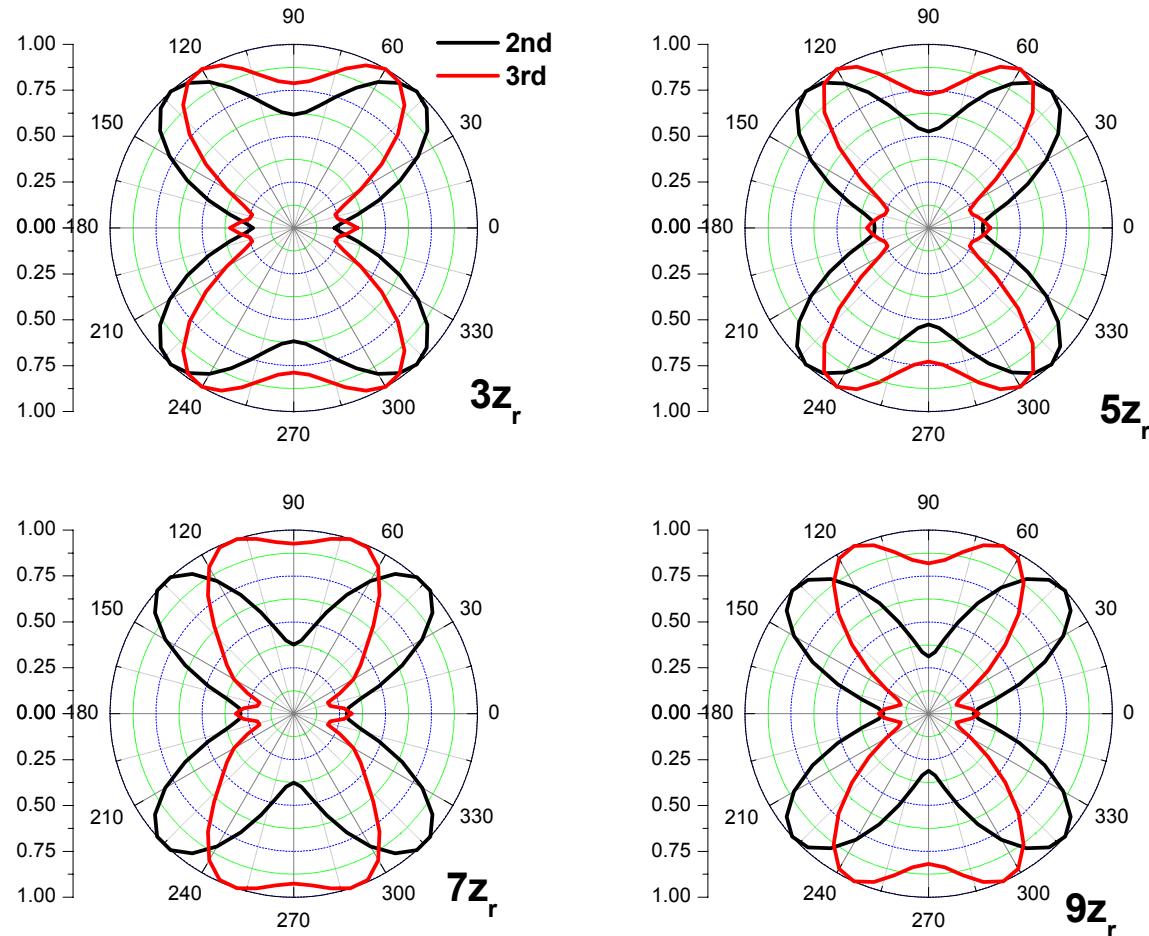
# Z-dependence: Spectra



# Target thickness dependence



# Target thickness: angular distribution



# Conclusion



Nonlinear Thomson scattering emission spectra and angular distribution are sensitive to the laser intensity, pulse shape and target configuration both on the microscopic and macroscopic level, showing spectral shift, broadening and mixing between harmonics.

# References



1. E. S. Sarachik and G.T. Schappert, "Classical theory of the scattering of intense laser radiation by free electrons," Phys. Rev. D **1**, 2738-2753 (1970).
2. S. K. Ride, E. Esarey, and M. Baine, "Nonlinear Thomson scattering of intense laser pulses from beams and plasmas," Phys. Rev. E **48**, 3003-3021 (1993).
3. S-Y Chen, A. Maksimchuk, and D. Umstadter, "Experimental observation of relativistic nonlinear Thomson scattering," Nature **396**, 653-655 (1998).
4. K. Ta Phus et al., "Synchrotron-like radiation produced from an intense femtosecond laser system," OSA Conference on Applications of High Field and Short Wavelength Sources IX, October 21–24, 2001, Palm Springs, California.
5. B. Quesnel and P. Mora, "Theory and simulation of the interaction of ultraintense laser pulses with electrons in vacuum," Phys. Rev. E **58**, 3719-3732 (1998), and references therein.
6. J. D. Jackson, *Classical Electrodynamics* (Wiley, New York, 1976).
7. Y. Ueshima, Y. Kishimoto, A. Sasaki and T. Tajima, "Laser Larmor X-ray radiation from low-Z matter," Laser Part. Beams **17**, 45-58 (1999).